



Reciprocating Internal Combustion Engines

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2012 Princeton-CEFRC
Summer Program on Combustion
Course Length: 9 hrs
(Wed., Thur., Fri., June 27-29)

Hour 4

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Hour 4: Atomization, Drop Breakup/Coalescence



Short course outline:

Engine fundamentals and performance metrics, computer modeling supported by in-depth understanding of fundamental engine processes and detailed experiments in engine design optimization.

Day 1 (Engine fundamentals)

Hour 1: IC Engine Review, 0, 1 and 3-D modeling

Hour 2: Turbochargers, Engine Performance Metrics

Hour 3: Chemical Kinetics, HCCI & SI Combustion

Day 2 (Spray combustion modeling)

Hour 4: Atomization, Drop Breakup/Coalescence

Hour 5: Drop Drag/Wall Impinge/Vaporization

Hour 6: Heat transfer, NOx and Soot Emissions

Day 3 (Applications)

Hour 7: Diesel combustion and SI knock modeling

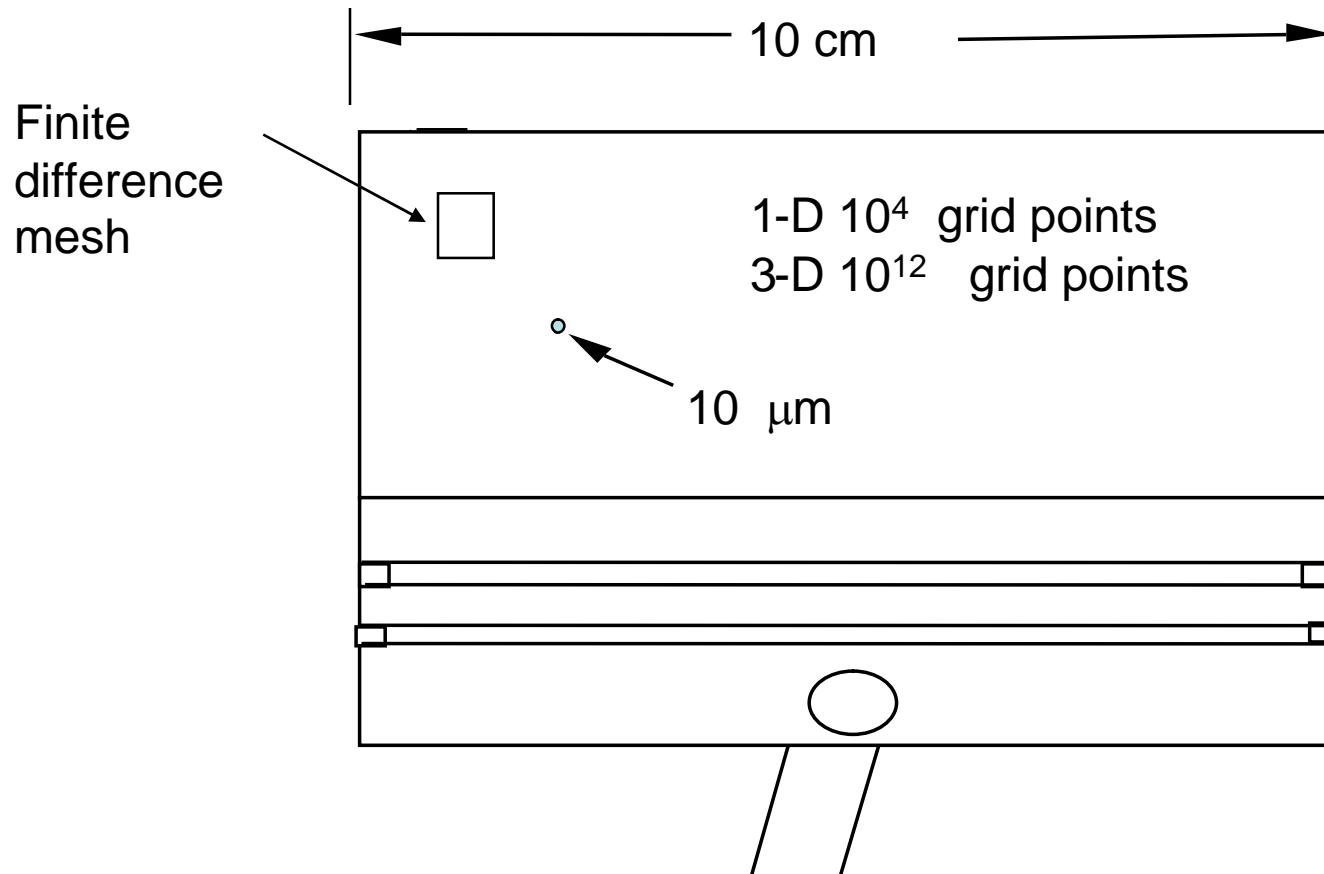
Hour 8: Optimization and Low Temperature Combustion

Hour 9: Automotive applications and the Future





Resolution – predictive models



Models will not be entirely predictive over next decade
Accurate submodels will be needed for detailed spray processes
(e.g., drop drag, drop turbulence interaction, vaporization, atomization,
drop breakup, collision and coalescence, and spray/wall interaction)





Amsden et al. 1997

Governing Equations

Gas phase
Liquid phase
Turbulence

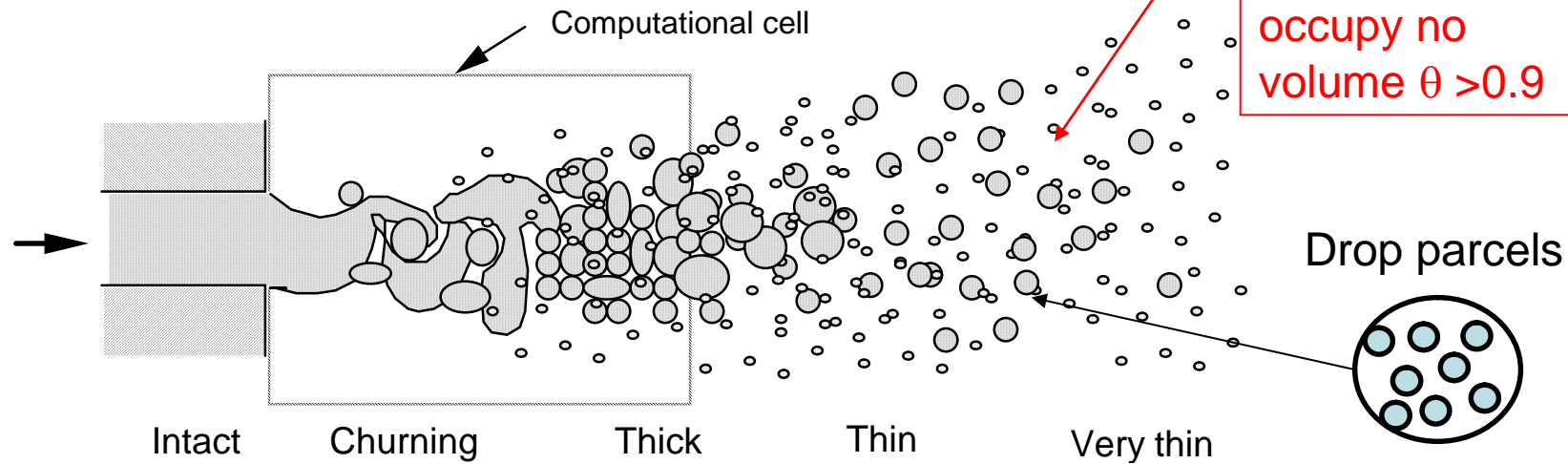
$$\bullet \rightarrow f = f(\mathbf{x}, \mathbf{v}, r, T_d; t)$$

$\mathbf{x}, \mathbf{v}, r, T_d$

Gas void fraction and drop number density

$$\theta = 1 - \int_{Vol} \left(\iiint \frac{4}{3} \pi r^3 f dr d\mathbf{v} dT_d \right) dVol / Vol$$

Two-Phase Flow Regimes





Spray Modeling

Dukowicz 1980

- Concept of using “drop parcels”

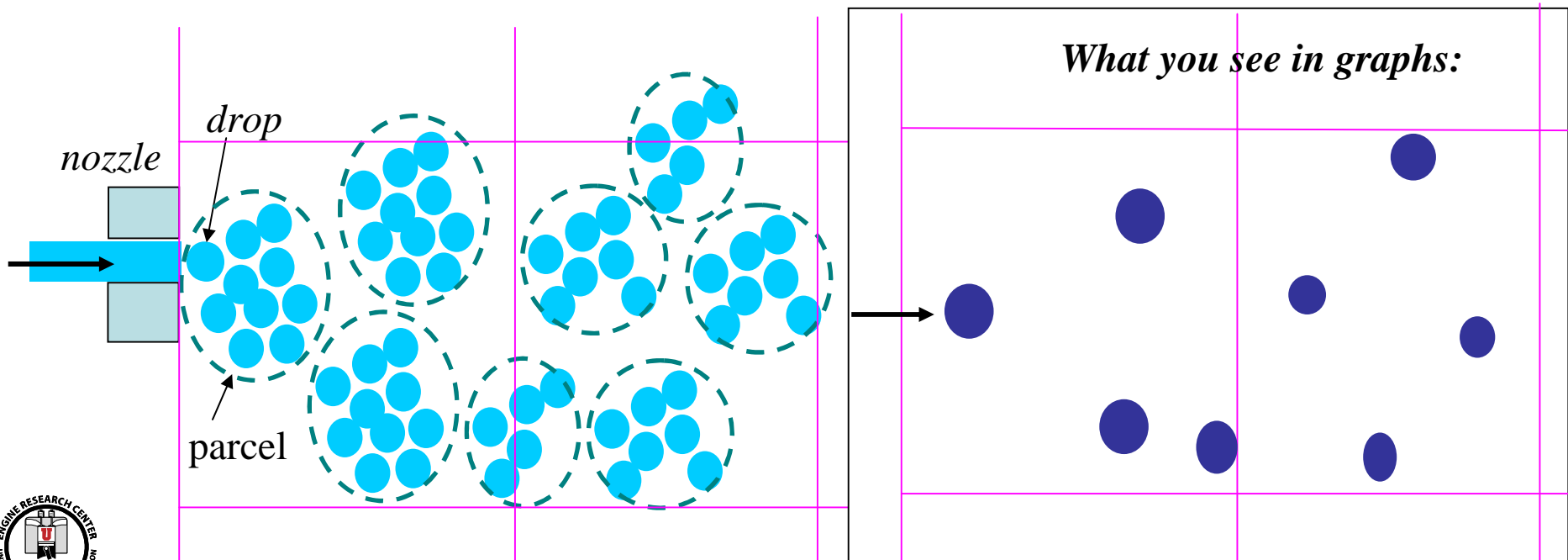
For typical heavy-duty diesel, injected fuel per cycle (75% load): 0.160 g

One spray plume: $m_{fuel} = 0.160/6 = 0.0267$ g

If average SMD = $10 \mu\text{m} \rightarrow m_{drop} = 3.8 \times 10^{-10}$ g

of drops in the domain = $0.0267 \text{g} / m_{drop} = 7.1 \times 10^7$

Impractical to track individual fuel drops – group identical drops into ‘parcels’





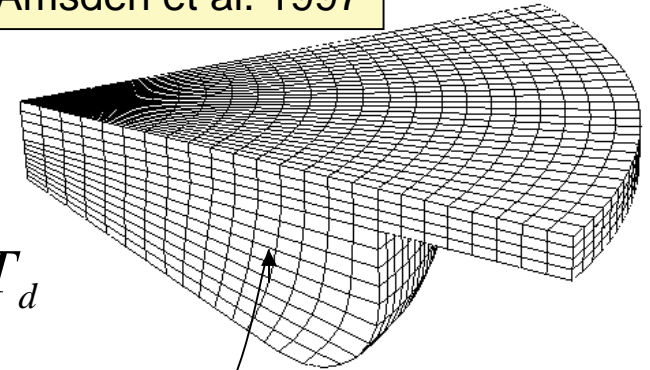
Eulerian Gas Phase

Amsden et al. 1997

Mass conservation (species)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = - \underbrace{\iiint \rho_l 4\pi r^2 R f dr d\mathbf{v} dT_d}_{\text{Vapor source}}$$

$R = dr/dt$ - Vapor source



Momentum conservation

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = -\nabla p - \underbrace{\nabla \left(\frac{2}{3} \rho k \right) + \nabla \boldsymbol{\tau}}_{\text{Turbulent and viscous stress}} + \underbrace{F^s}_{\text{Rate of momentum gain due to spray - drop drag}} + \rho \mathbf{g}$$

Turbulent and viscous stress

Rate of momentum gain due to spray – drop drag





Gas Phase (2)

Amsden et al. 1989

Internal energy conservation

$$\frac{\partial \rho I}{\partial t} + \nabla \cdot (\rho \mathbf{u} I) = -P \nabla \cdot \mathbf{u} - \nabla \cdot \mathbf{J} + \underbrace{\rho \varepsilon}_{\text{Turbulence dissipation}} + \underbrace{\dot{Q}^c}_{\text{Combustion heat release}} + \underbrace{\dot{Q}^s}_{\text{Energy due to Spray - vaporization}}$$

Turbulence
dissipation

Energy due to
Spray - vaporization

Heat flux

$$\mathbf{J} = -\lambda \nabla T - \rho D \sum_m h_m \nabla (\rho_m / \rho)$$

Equations of state

$$p = RT \sum_m \rho_m / W_m$$

Specific heat, enthalpy from JANAF data





Liquid Phase

Amsden, 1997

Spray drop number conservation $f = f(\mathbf{x}, \mathbf{v}, r, T_d, y, \dot{y}; \mathbf{t})$

$$\frac{\partial f}{\partial t} + \underbrace{\nabla_{\mathbf{x}} \cdot (f\mathbf{v}) + \nabla_{\mathbf{v}} \cdot (f\mathbf{F})}_{\substack{\mathbf{F} = d\mathbf{v}/dt \\ \text{drop drag}}} + \underbrace{\frac{\partial}{\partial r}(fR) + \frac{\partial}{\partial T_d}(f\dot{T}_d)}_{\substack{R = dr/dt \\ \text{Vaporization and heating}}} + \underbrace{\frac{\partial}{\partial y}(f\dot{y}) + \frac{\partial}{\partial \dot{y}}(f\ddot{y})}_{\substack{\text{Drop distortion}}} = \underbrace{\dot{f}_{coll} + \dot{f}_{bu}}_{\substack{\text{Drop breakup,} \\ \text{coalescence}}}$$

Spray exchange functions

$$\mathbf{F}^s = - \int f \rho_d (4/3 \pi r^3 \mathbf{F}' + 4\pi r^2 R \mathbf{v}) d\mathbf{v} dr dT_d dy d\dot{y}$$

$$\dot{Q}^s = - \int f \rho_d \left\{ 4\pi r^2 R \left[I_l + \frac{1}{2}(\mathbf{v} - \mathbf{u})^2 \right] + 4/3 \pi r^3 [c_l \dot{T}_d + \mathbf{F}' \cdot (\mathbf{v} - \mathbf{u} - \mathbf{u}')] \right\} d\mathbf{v} dr dT_d dy d\dot{y}$$

Work done by drop drag forces $\dot{W}^s = - \int f \rho_d 4/3 \pi r^3 \mathbf{F}' \cdot \mathbf{u}' d\mathbf{v} dr dT_d dy d\dot{y}$





Lagrangian drop - liquid phase

Amsden, 1997

Discrete Drop Model

drop position

$$\frac{d\mathbf{x}}{dt} = \mathbf{v}$$

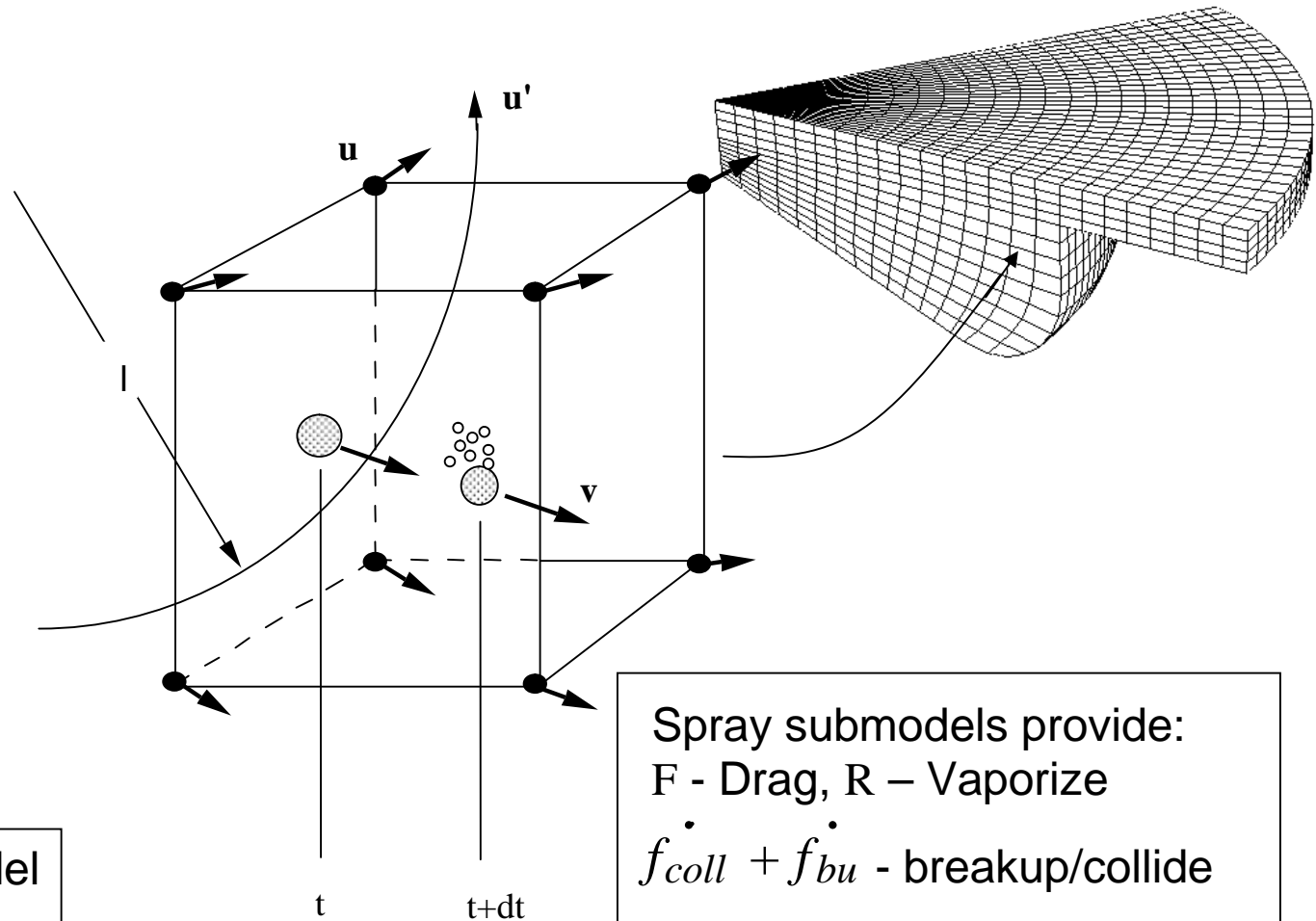
drop velocity

$$\frac{d\mathbf{v}}{dt} = \mathbf{F}$$

drop size

$$\frac{dr}{dt} = \mathbf{R}$$

Turbulence model provides: l, \mathbf{u}'



Spray submodels provide:
 \mathbf{F} - Drag, \mathbf{R} - Vaporize
 $\dot{f}_{coll} + \dot{f}_{bu}$ - breakup/collide
 Initial data:
 \mathbf{v}, r, T_d - Atomization model





Amsden, 1997

Turbulence Model (RANS)

Kinetic energy

$$\frac{\partial \rho k}{\partial t} + \nabla \cdot (\rho \mathbf{u} k) = -\frac{2}{3} \rho k \nabla \cdot \mathbf{u} + \underbrace{\boldsymbol{\tau} : \mathbf{u}}_{\text{Production due to mean flow}} + \nabla \cdot \left[\left(\frac{\mu}{Pr_k} \right) \nabla k \right] \underbrace{-\rho \varepsilon + \dot{W}^s}_{\text{Rate of work to disperse drops}}$$

Dissipation

Production due to mean flow

Rate of work to disperse drops

Dissipation rate

$$\frac{\partial \rho \varepsilon}{\partial t} + \nabla \cdot (\rho \mathbf{u} \varepsilon) = - \left(\frac{2}{3} C_{\varepsilon 1} - C_{\varepsilon 3} \right) \rho \varepsilon \nabla \cdot \mathbf{u} + \nabla \cdot \left[\left(\frac{\mu}{Pr_\varepsilon} \right) \nabla \varepsilon \right] + \frac{\varepsilon}{k} \left[C_{\varepsilon 1} \boldsymbol{\tau} : \nabla \mathbf{u} - C_{\varepsilon 2} \rho \varepsilon + C_s \dot{W}^s \right]$$

Turbulence diffusivity

$$D = C_\mu k^2 / \varepsilon$$

Eddy size

$$l = C k^{3/2} / \varepsilon$$

Turbulence intensity

$$u'^2 = (2 k/3)$$



UW-ERC Multidimensional CFD models



<u>Submodel</u>	<u>Los Alamos</u>	<u>UW-Updated</u>	<u>References</u>
intake flow	assumed initial flow	compute intake flow	SAE 951200
heat transfer	law-of-the-wall	compressible, unsteady	SAE 960633
turbulence	standard k- ϵ	RNG k- ϵ /LES	CST 106, 1995
nozzle flow	none	cavitation modeling	SAE 1999-01-0912
atomization	Taylor Analogy	surface-wave-growth Kelvin Hemholtz Rayleigh Taylor	SAE 960633 SAE 980131 CST 171, 1998
drop breakup	Taylor Analogy	Rayleigh Taylor	Atom. Sprays 1996
drop drag	rigid sphere	drop distortion	SAE 960861
wall impinge	none	rebound-slide model wall film/splash	SAE 880107 SAE 982584
collision/coalesce	O'Rourke	shattering collisions	Atom. Sprays 1999
vaporization	single component low pressure	multicomponent fuels high pressure	SAE 2000-01-0269 SAE 2001-01-0998
ignition	Arrhenius	reduced chemistry	SAE 2004-01-0558
combustion	Arrhenius	CTC/GAMUT reduced kinetics	SAE 2004-01-0102 SAE 2003-01-1087
NOx	Zeldo'vich	Extended Zeldo'vich	SAE 940523
soot	none	Hiroyasu & Surovkin Nagle Strickland oxidation	SAE 960633 SAE 980549

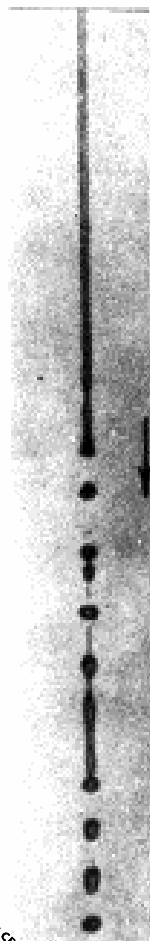




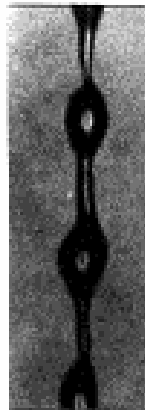
Review of atomization models (Single Hole Nozzle)

Reitz & Bracco, 1982

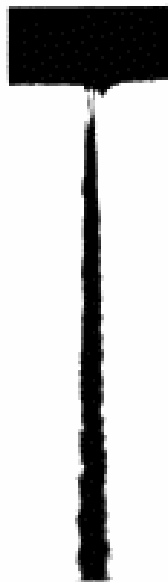
Four main jet breakup regimes:
Rayleigh, first wind-induced, second wind-induced and atomization



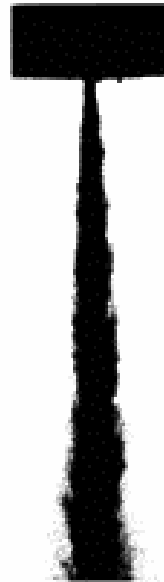
(a)



(b)



(c)



(d)

- a.) Rayleigh breakup
Drop diameters > jet diameter.
Breakup far downstream nozzle
- b.) First wind-induced regime
Drop diameter ~ jet diameter.
Breakup far downstream of nozzle
- c.) Second wind-induced regime
Drop sizes < jet diameter.
Breakup starts close to nozzle exit
- d.) Atomization regime
Drop sizes << jet diameter.
Breakup at nozzle exit.

Growth of small disturbances
initiates liquid breakup





Lee & Reitz, 2010

Nozzle flow - cavitation

Homogeneous Equilibrium Model

- single phase mixture of vapor and liquid
- considers variable compressibility of mixture.

(1) Sonic Speed of mixture

: function of void fraction

$$\alpha = \frac{\rho_l - \rho}{\rho_l - \rho_v} \quad \begin{array}{l} \alpha=0 \text{ for pure liquid} \\ \alpha=1 \text{ for pure vapor} \end{array}$$

$$\frac{1}{a^2} = [\alpha\rho_v + (1-\alpha)\rho_l] \left[\frac{\alpha}{\rho_v a_v^2} + \frac{1-\alpha}{\rho_l a_l^2} \right]$$

(Wallis, 1967)

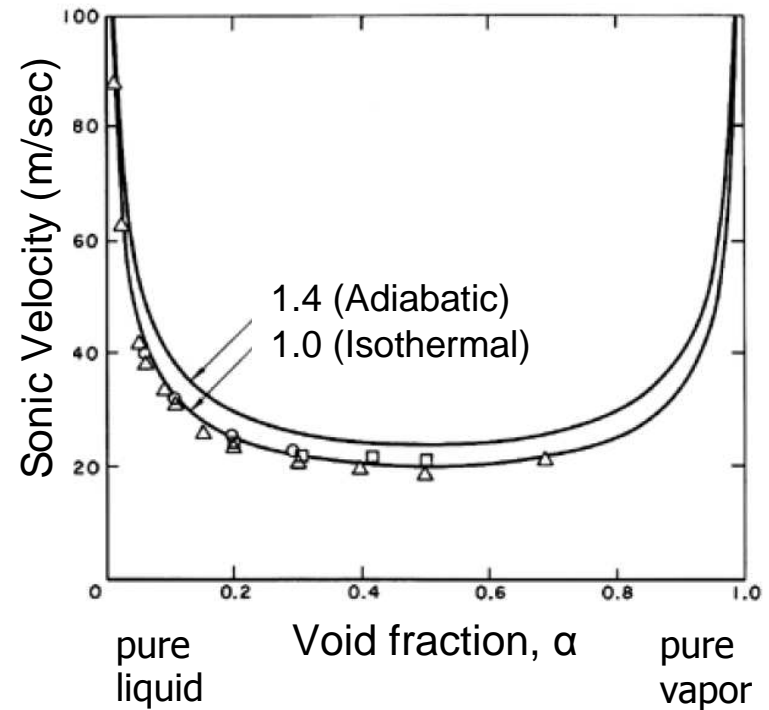
(2) Equation of State of mixture

: by integrating $dP = a^2 d\rho$ (Schmidt, 1997)

$$P = P_l^{sat} + P_{vl} \log \left[\frac{\rho_v a_v^2 (\rho_l + \alpha(\rho_v - \rho_l))}{\rho_l (\rho_v a_v^2 - \alpha(\rho_v a_v^2 - \rho_l a_l^2))} \right]$$

$$P_{vl} = \frac{\rho_v a_v^2 \rho_l a_l^2 (\rho_v - \rho_l)}{\rho_v^2 a_v^2 - \rho_l^2 a_l^2}$$

$$P_l^{sat} = P_v^{sat} + P_{vl} \log \left[\frac{\rho_v^2 a_v^2}{\rho_l^2 a_l^2} \right]$$



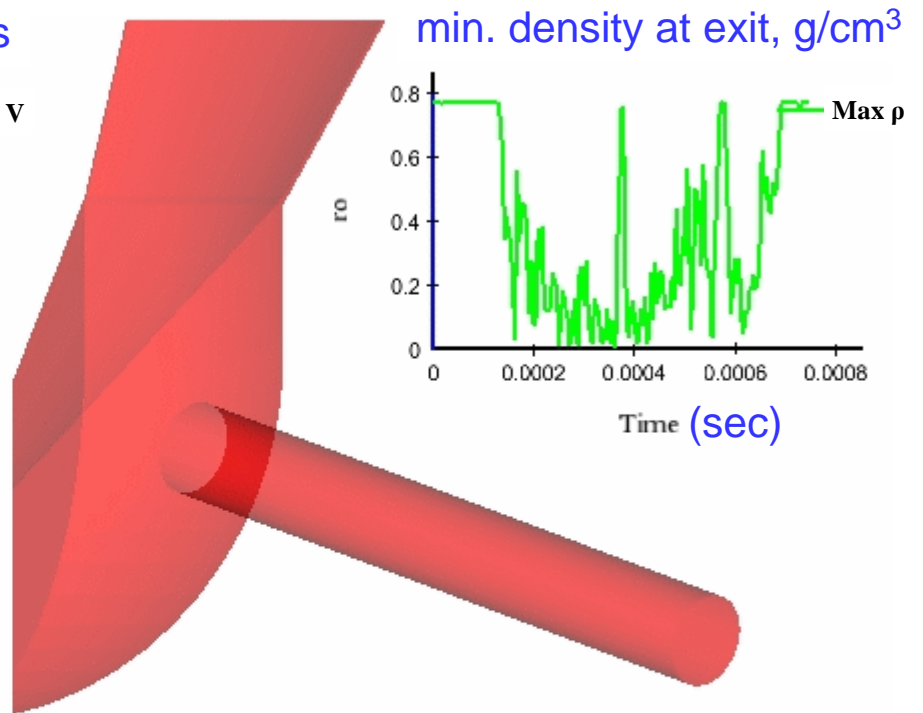
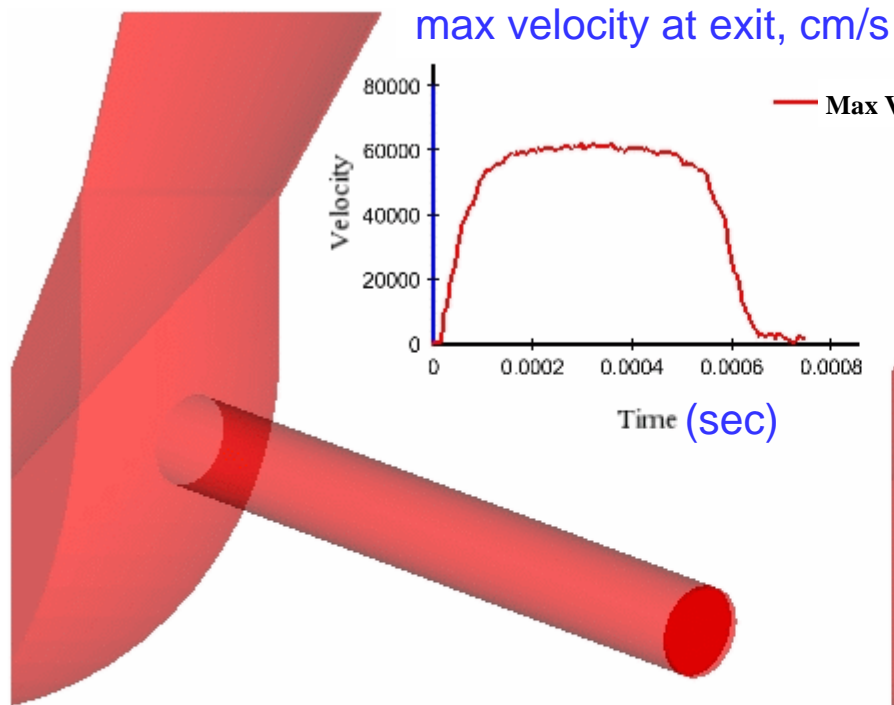
Sonic velocity in bubbly air/water mixture at atmospheric pressure
Brennen (1995)



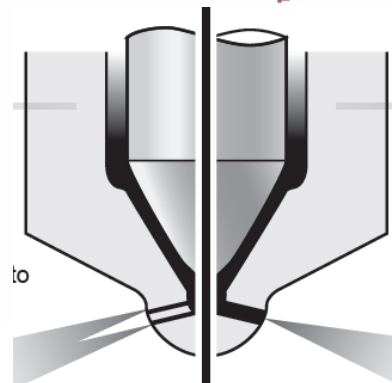


Nozzle flow - cavitation

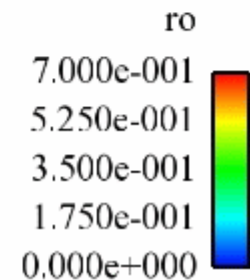
Lee & Reitz, 2010



streamline and exit velocity



density and iso-surface (ρ=0.35g/cm³)

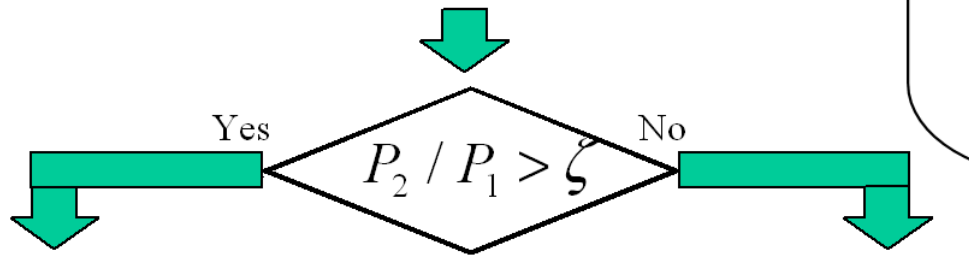




Cavitation Inception

Account for effects of nozzle geometry

Cavitation if $P < P_V$



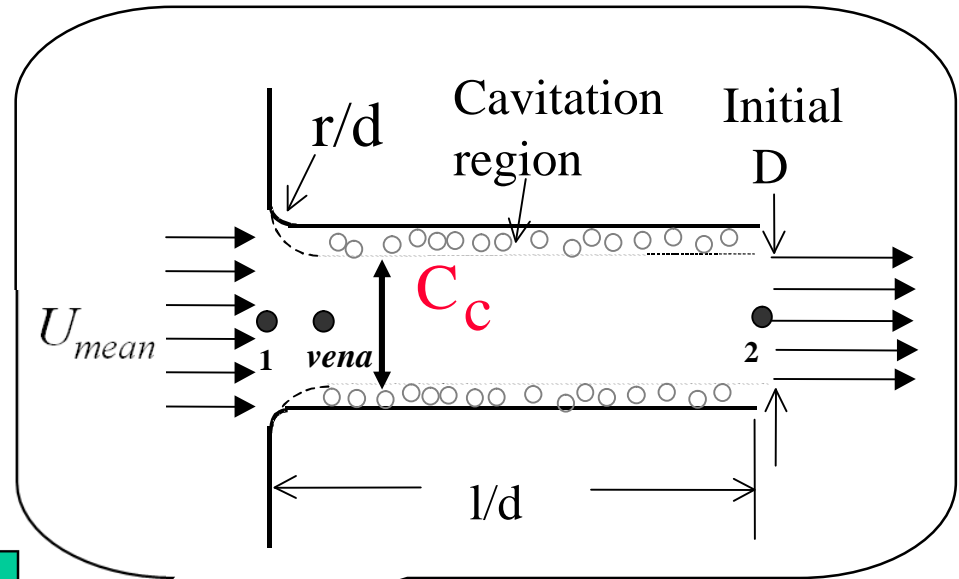
Cavitating flow

$$\zeta = \frac{1}{2(C_c - C_c^2)}$$

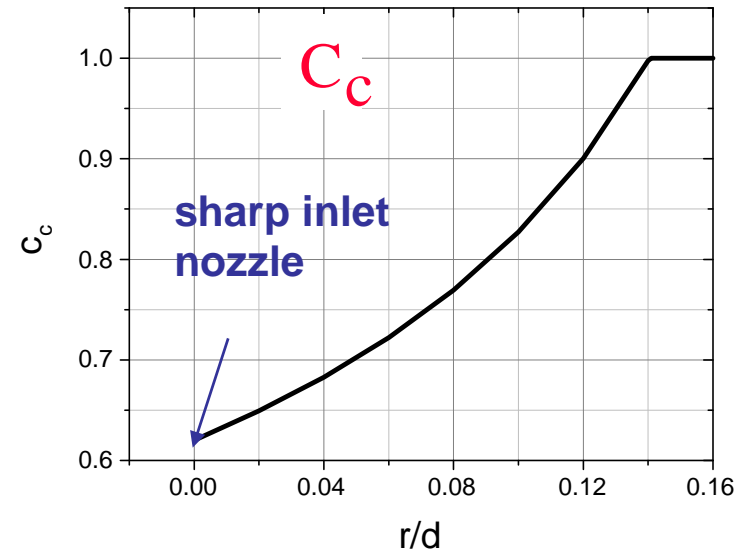
Non-cavitating flow

Contraction coefficient (Nurick (1976))

$$C_c = \left[\left(\frac{1}{0.62} \right)^2 - 11.4r/d \right]^{-1/2}$$



Sarre SAE 1999-01-0912





ERC Nozzle Flow Model

Cavitating flow

Yes

No

Non-cavitating flow

$$P_2 / P_1 > \zeta$$

Nozzle discharge coefficient

$$C_d = C_c \sqrt{\frac{P_1 - P_v}{P_1 - P_2}}$$

Effective injection velocity

$$u_{eff} = \frac{2C_c P_1 - P_2 + (1 - 2C_c) P_v}{C_c \sqrt{2\rho(P_1 - P_v)}}$$

Effective nozzle area

$$A_{eff} = \frac{2C_c^2 (P_1 - P_v)}{2C_c P_1 - P_2 + (1 - 2C_c) P_v} A$$

Nozzle discharge coefficient

Lichtarowicz (1965)

$$C_d = 0.827 - 0.0085 l/d$$

Effective injection velocity

$$u_{eff} = C_d \sqrt{\frac{2(P_1 - P_2)}{\rho}}$$

Effective nozzle area

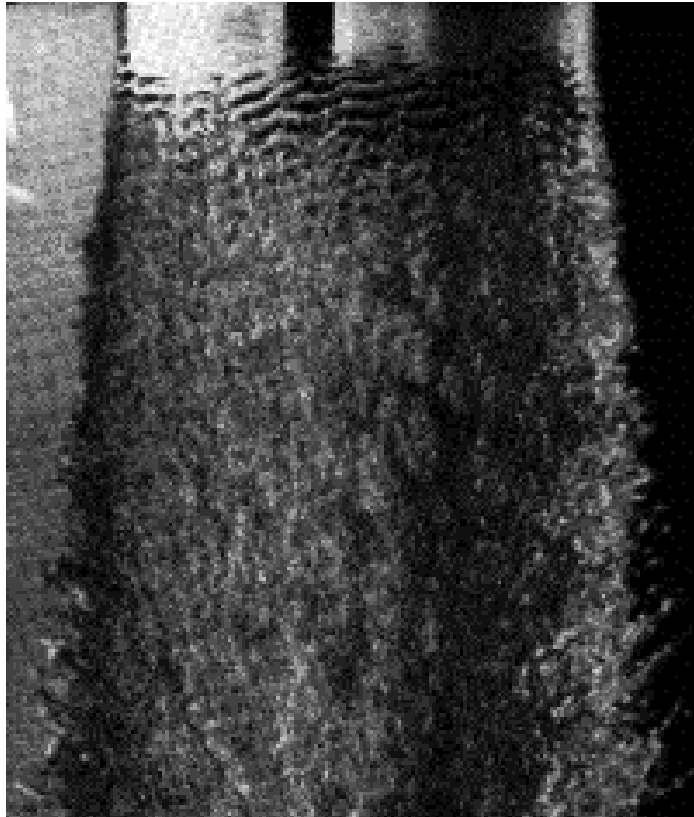
$$A_{eff} = A$$

Sarre SAE 1999-01-0912





Atomization - "Wave" breakup model

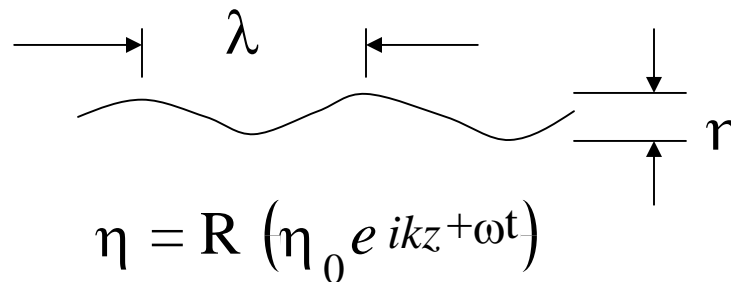


Taylor & Hoyt, 1983

High speed photograph of water jet close to nozzle exit (at top) in the second wind-induced breakup regime showing surface wave instability growth and breakup

Reitz & Bracco, 1982

Kelvin Helmholtz Jet Breakup Model



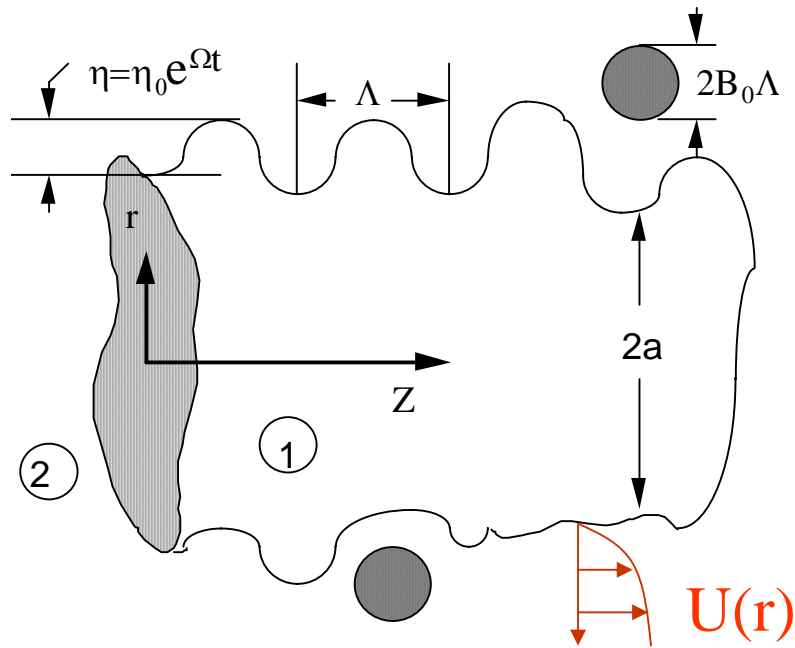
Linear Stability Theory:

Cylindrical liquid jet issuing from a circular orifice into a stationary, incompressible gas.

Relate growth rate, ω , of perturbation to wavelength $\lambda = 2\pi/k$



Hour 4: Atomization, Drop Breakup/Coalescence



Linearized analysis

Reitz & Bracco, 1982

U = Jet velocity

Surface waves breakup on jet or "blob"

$$\eta = R (\eta_0 e^{ikz + \omega t})$$

Equation of liquid surface: $r = a + \eta$,

Axisymmetric fluctuating pressure, axial velocity, and radial velocity for both liquid and gas phases.

Fluctuations described by continuity equation

$$\frac{\partial u_i}{\partial z} + \frac{1}{r} \frac{\partial}{\partial r} (r v_i) = 0$$

plus linearized equations of motion for the liquid and the gas,

$$\text{Axial: } \frac{\partial u_i}{\partial t} + U(r) \frac{\partial u_i}{\partial z} + v_i \frac{dU}{dr} = - \frac{1}{\rho} \frac{\partial p}{\partial z} + \frac{\mu}{\rho} \left[\frac{\partial^2 u_i}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_i}{\partial r} \right) \right]$$





Analysis (Cont.)

$$\text{Radial: } \frac{\partial v_i}{\partial t} + U_i(r) \frac{\partial v_i}{\partial z} = -\frac{1}{\rho_i} \frac{\partial p_i}{\partial r} + \frac{\mu_i}{\rho_i} \left[\frac{\partial^2 v_i}{\partial z^2} + \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial r v_i}{\partial r} \right) \right]$$

Gas is assumed to be inviscid $U(r) = U$ - slip

With $\eta \ll a$, the gas equations give the pressure at the interface $r = a$

$$p_2 = -\rho_2 \left(U - i \frac{\omega}{k} \right)^2 k \eta \frac{K_0(ka)}{K_1(ka)}$$

Boundary conditions-

Kinematic, tangential and normal stress at the interface:

$$v_1 = \mathbf{w} = \frac{\partial \eta}{\partial t}, \quad \frac{\partial u_1}{\partial r} = -\frac{\partial v_1}{\partial z}$$

$$-p_1 + 2v_1\rho_1 \frac{\partial v_1}{\partial r} - \frac{\sigma}{a^2} \left(\eta + a^2 \frac{\partial^2 \eta}{\partial z^2} \right) + p_2 = 0$$





Dispersion Relationship

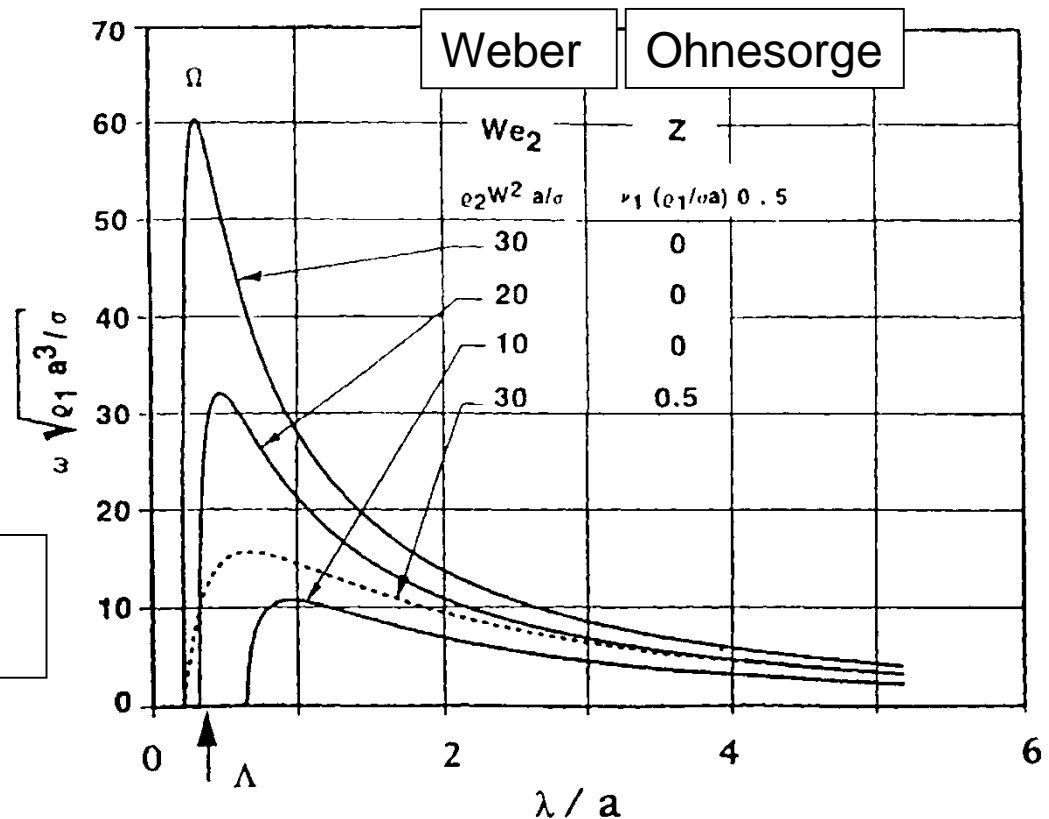
Reitz, 1988

$$\omega^2 + 2v_1 k^2 \omega \left[\frac{I_1'(ka)}{I_0(ka)} - \frac{2kl}{k^2 + l^2} \frac{I_1(ka)}{I_0(ka)} \frac{I_1'(la)}{I_0(la)} \right] = \frac{\sigma k}{\rho_1 a^2} (1 - k^2 a^2) \left(\frac{l^2 - k^2}{l^2 + k^2} \right) \frac{I_1(ka)}{I_0(ka)}$$

$$+ \frac{\rho_2}{\rho_1} (U - i\omega/k)^2 k^2 \left(\frac{l^2 - k^2}{l^2 + k^2} \right) \frac{I_1(ka)K_0(ka)}{I_0(ka)K_1(ka)}$$

Maximum wave growth rate characterizes fastest growing waves which are responsible for breakup (as a function of Weber and Ohnesorge numbers)

Maximum wave growth rate and length scale: Ω and Λ





Curvefit of Dispersion Equation

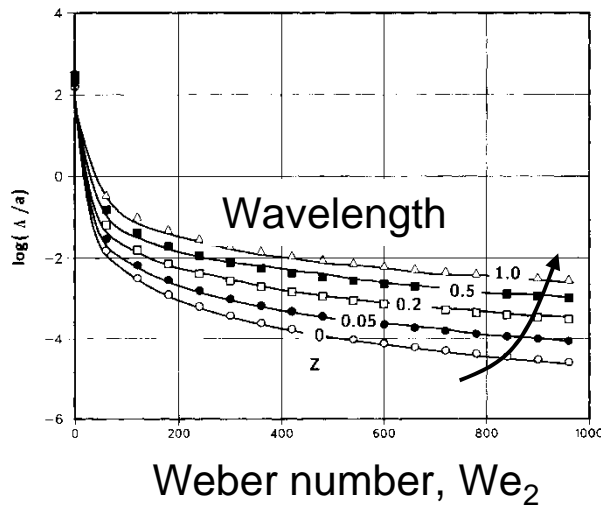
Reitz, 1988

$$\frac{\Delta}{a} = 9.02 \frac{(1 + 0.45 Z^{0.5})(1 + 0.4T^{0.7})}{(1 + 0.87 We_2^{1.67})^{0.6}}$$

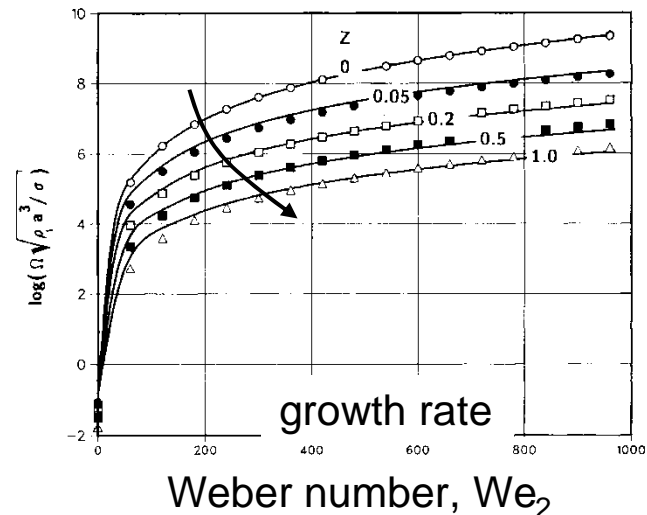
$$\Omega \left(\frac{\rho_1 a^3}{\sigma} \right)^{0.5} = \frac{0.34 + 0.38 We_2^{1.5}}{(1 + Z)(1 + 1.4T^{0.6})}$$

where $Z = \frac{We_1^{0.5}}{Re_1}$; $T = Z We_2^{0.5}$; $We_1 = \frac{\rho_1 U^2 a}{\sigma}$; $We_2 = \frac{\rho_2 U^2 a}{\sigma}$; $Re_1 = \frac{Ua}{\nu_1}$

Maximum growth rate increases and wavelength decreases with We
 Increased viscosity reduces growth rate and increases wave length



Ohnesorge number, Z



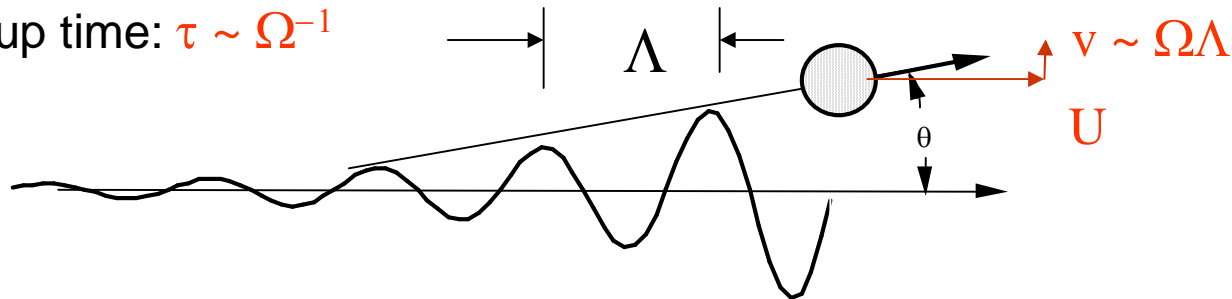


“Wave” atomization model

Reitz, 1988

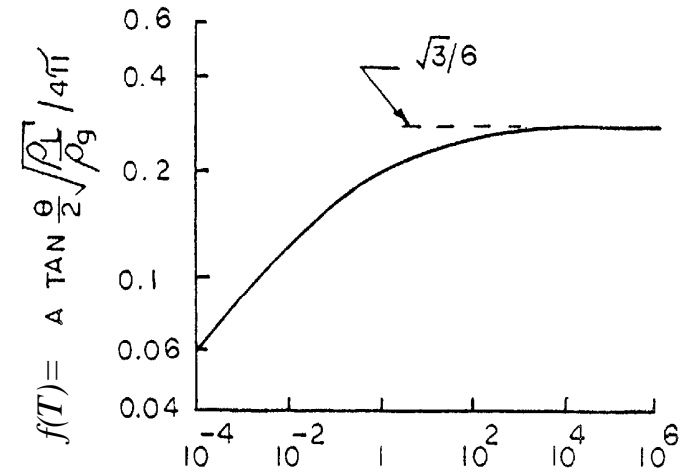
Drop size: $r = B\Lambda$

Breakup time: $\tau \sim \Omega^{-1}$



Spray angle prediction:

$$\tan \theta = \frac{v}{U} = \frac{1}{A} 4\pi \left(\frac{\rho_s}{\rho_l}\right)^{1/2} f(T)$$



Breakup length of the core (Taylor, 1940):

$$L = C a \sqrt{\frac{\rho_l}{\rho_2}} / f(T) \quad \text{where} \quad f(T) = \frac{\sqrt{3}}{6} [1 - \exp(-10T)]$$

$$T = \frac{\rho_L}{\rho_g} \left(\frac{Ra_L}{We_L} \right)^2$$



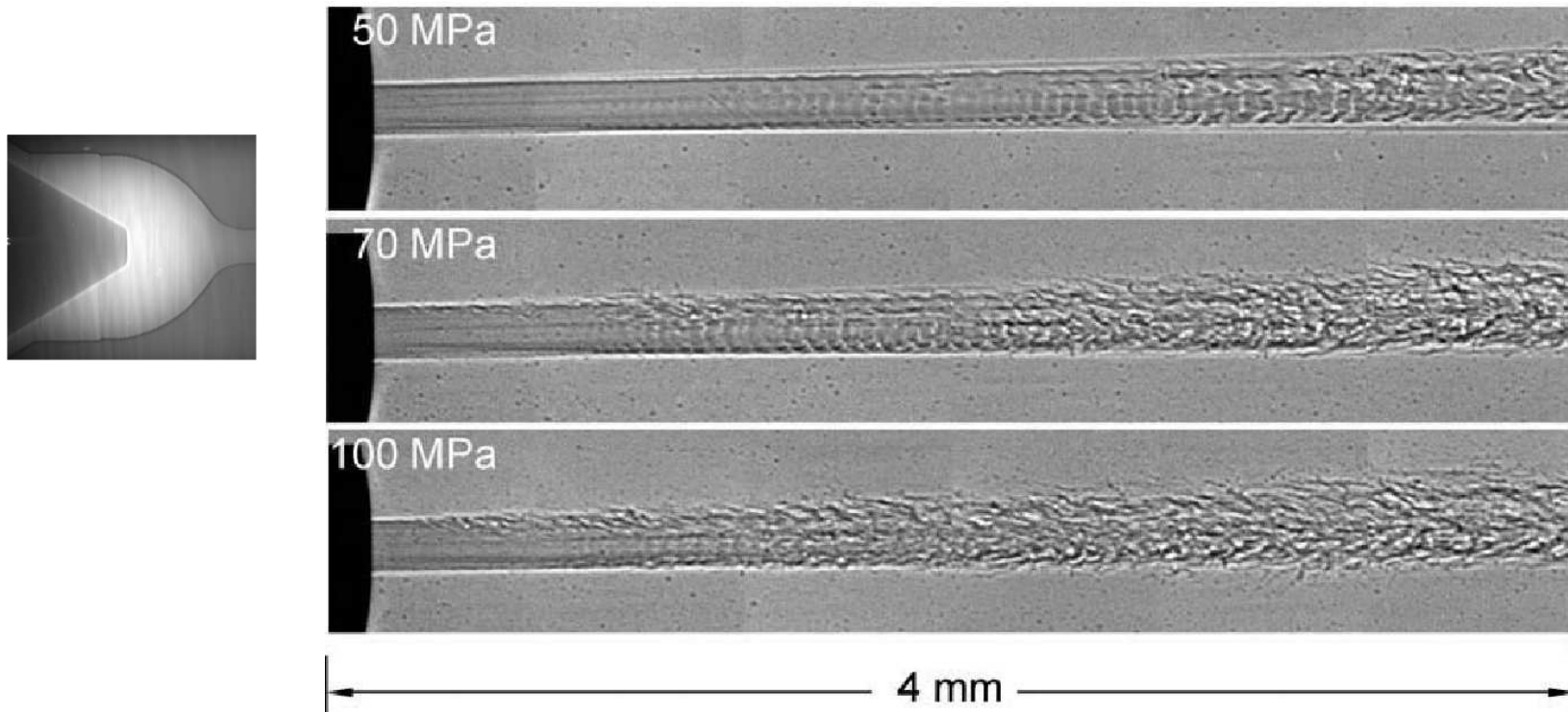


X-ray Phase-contrast imaging of high-pressure sprays

ANL Synchrotron-Based Ultrafast (150 ps) Single-Shot images

Surface instability waves produce ligaments

Breakup sensitive to injection pressure, fuel properties



(Hydroground nozzle, biodiesel, 1 ms injection duration in quasi-steady state)



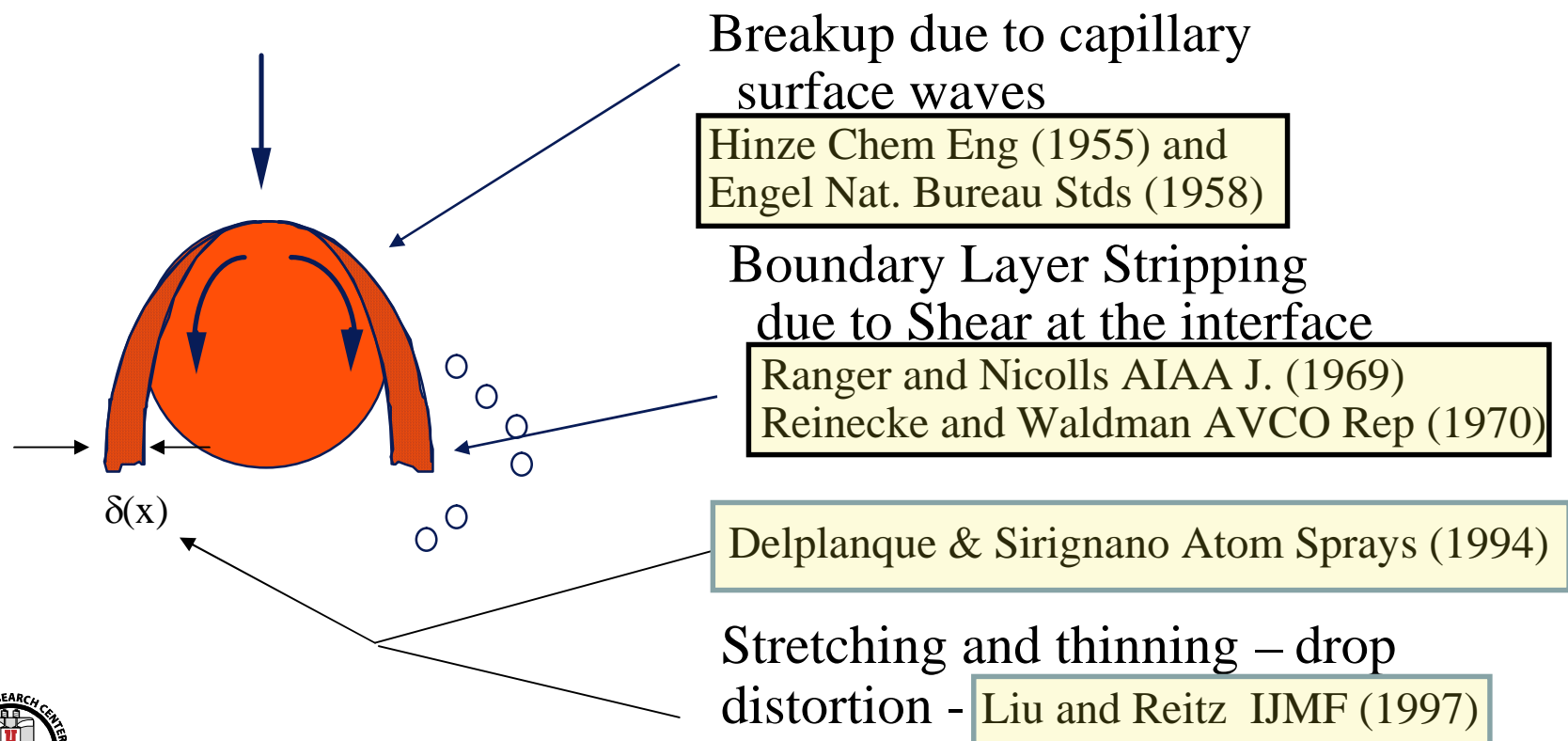
Gao, 2010



Drop breakup

Mechanisms of drop breakup at high velocities
poorly understood - Conflicting theories

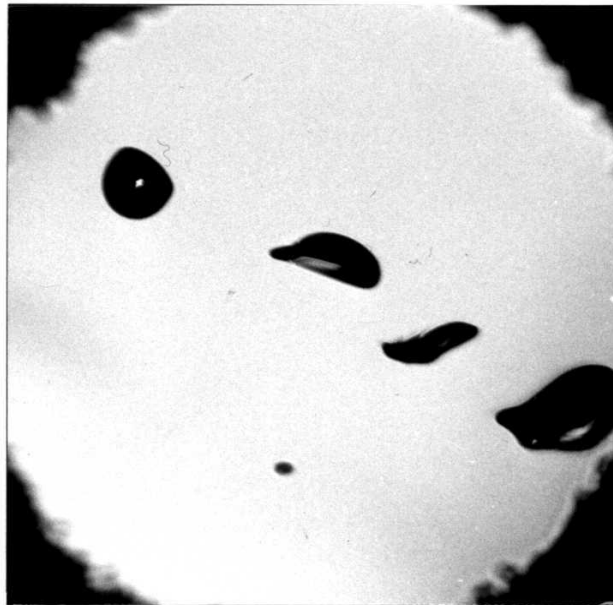
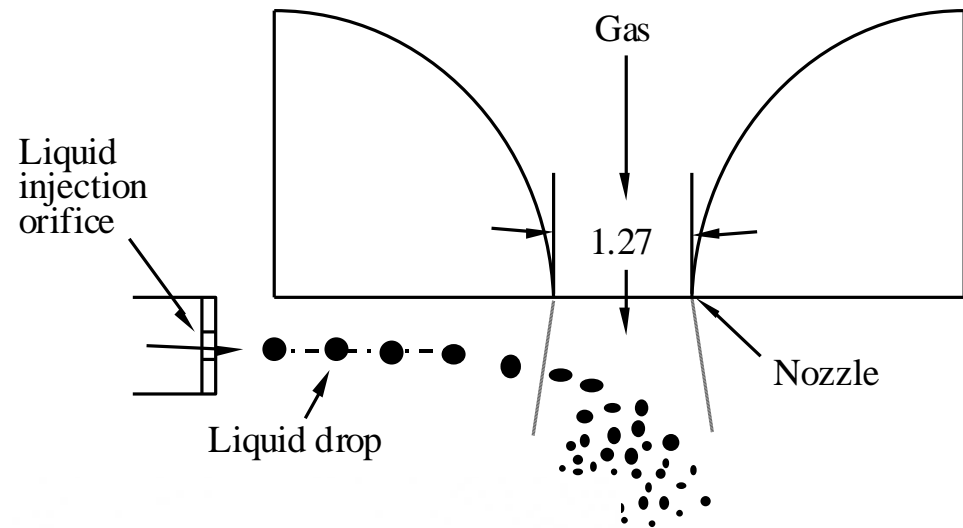
Bag, 'Shear' and 'Catastrophic' breakup regimes



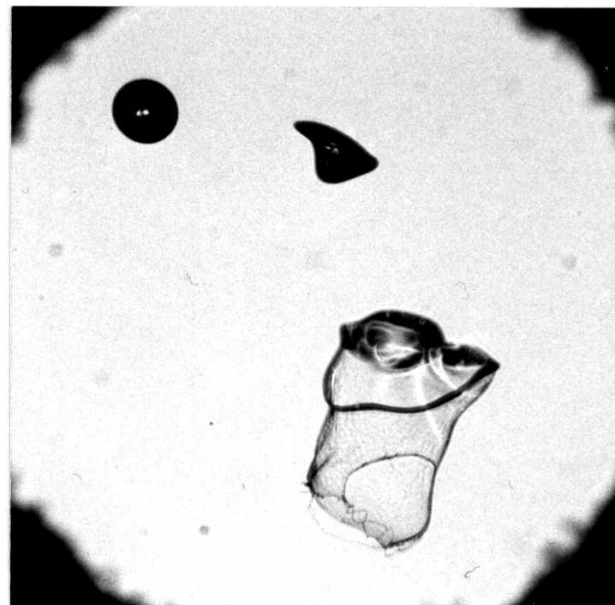


Low velocity drop breakup

Liu & Reitz, 1993



(a) bag breakup ($We=78$)



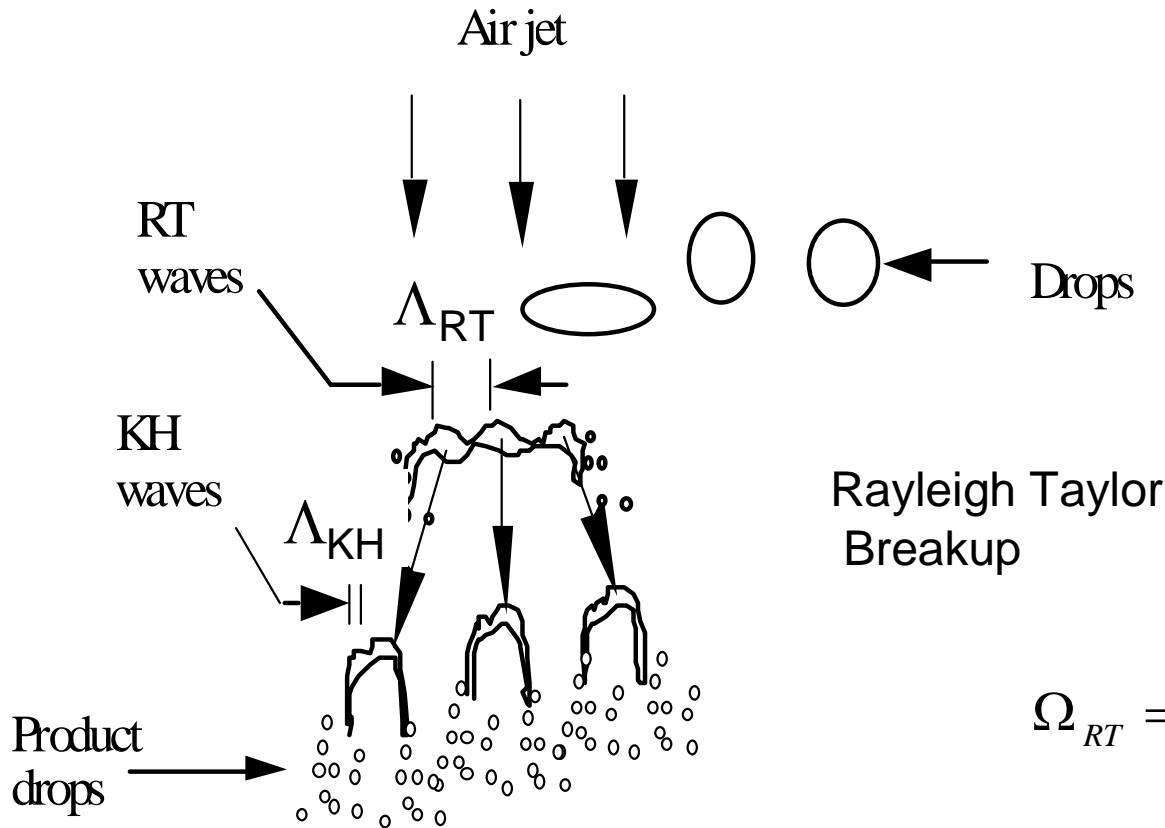
(b) bag breakup ($We=98$)



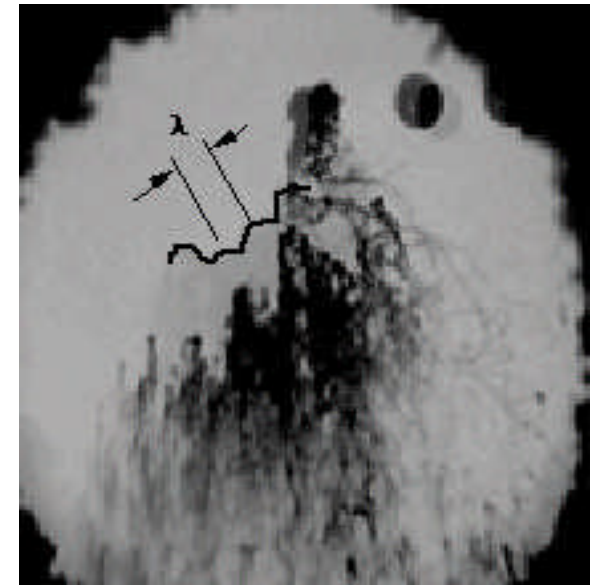


High speed drop breakup mechanism

Double pulse images



Rayleigh Taylor Breakup



$$\Omega_{RT} = \sqrt{\frac{2}{3\sqrt{\sigma}} \frac{[-g_t(\rho_l - \rho_g)]^{3/2}}{\rho_l + \rho_g}}$$

$g_t =$ acceleration

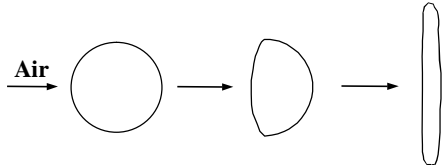
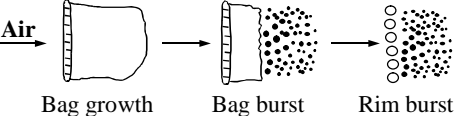
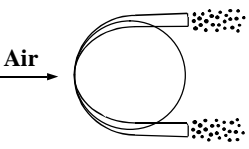
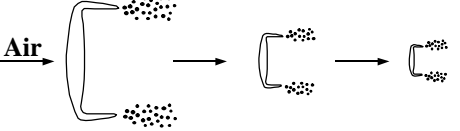
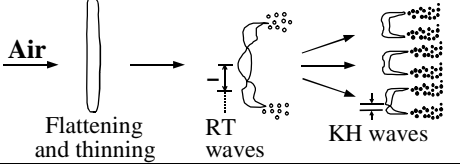
$$\Lambda_{RT} = \sqrt{\frac{-g_t(\rho_l - \rho_g)}{3\sigma}}$$

Hwang, 1996





Drop Breakup Regimes

Breakup stages	Deformation or breakup regimes	Breakup process	Weber number	References
First breakup stage	(1) Deformation and flattening		$We < 12$	
Second breakup stage	(b) Bag breakup		$12 \leq We \leq 100$ (including the Bag-and-Stamen breakup)	Pilch and Erdman
	(c) Shear breakup		$We < 80$	Ranger and Nicolls 1969
	(d) Stretching and thinning breakup		$100 \leq We \leq 350$	Liu and Reitz 1997
	(e) Catastrophic breakup		$350 \leq We$	Hwang et al. 1996



Lee & Reitz, 2001

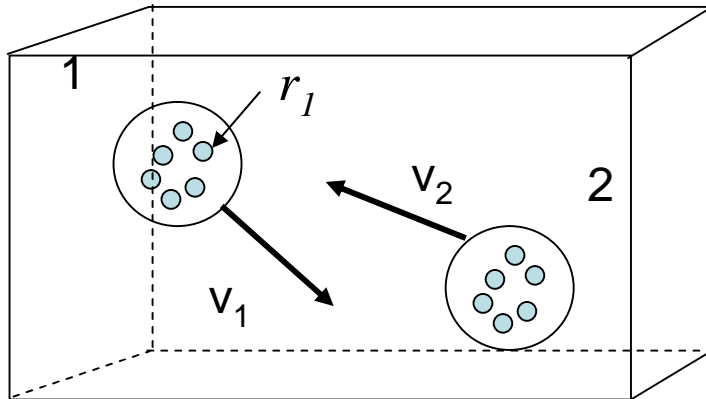
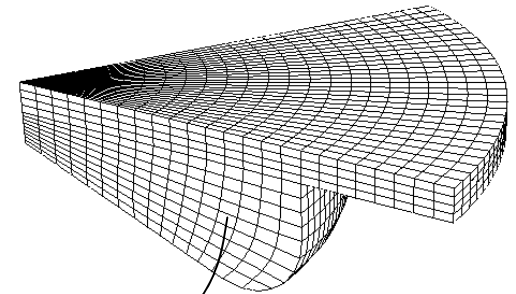


Drop collision modeling

O'Rourke PhD thesis 1981

Collision frequency

$$v_{12} = N_2 \pi (r_1 + r_2)^2 E_{12} |\mathbf{v}_1 - \mathbf{v}_2| / Vol$$



Number of collisions from Poisson process

Collision efficiency

$$E_{12} = \left(\frac{K}{K + 1/2} \right)^2 \sim 1$$

$$p(n) = e^{-v_{12}\Delta t} (v_{12}\Delta t)^n / n!$$

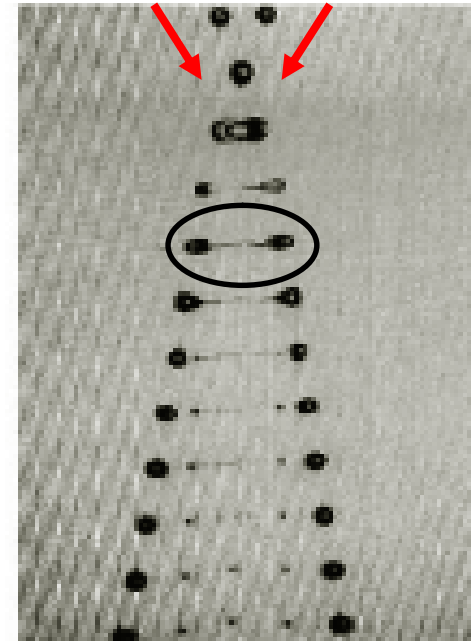
$$0 < p < 1 \text{ random number}$$

$$K = \frac{2 \rho_l |\mathbf{v}_1 - \mathbf{v}_2| r_2^2}{9 \mu_g r_1}$$



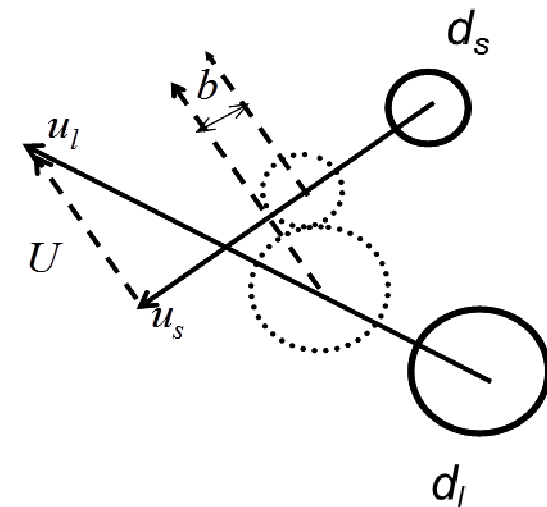
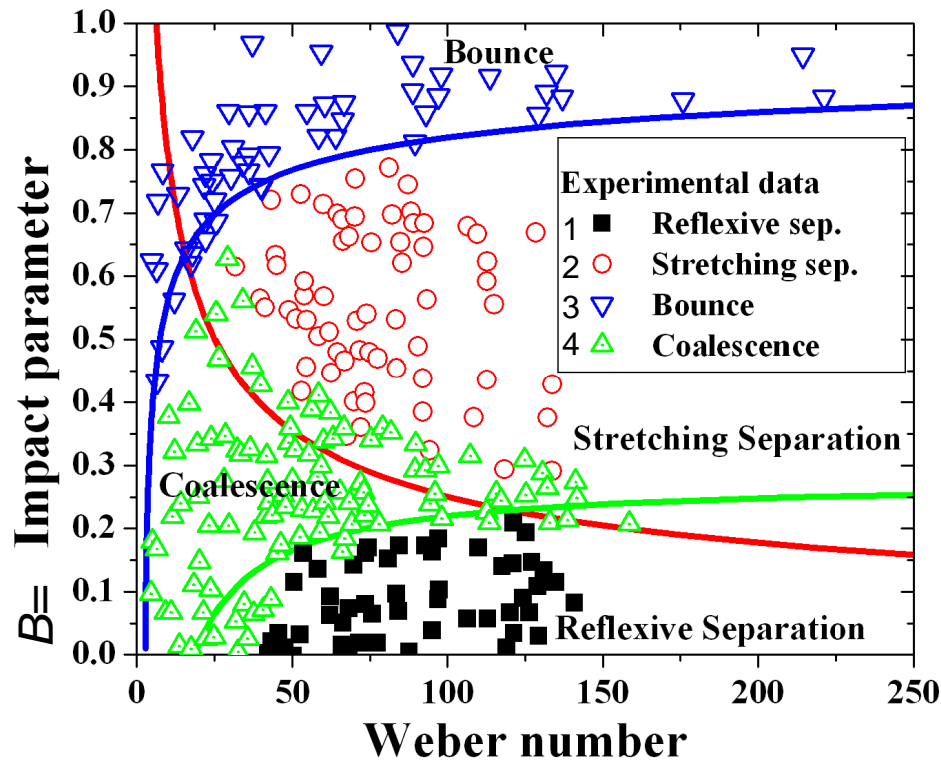


Munnannur, 2007



Drop collision and coalescence

1. Reflexive vs. surface energy
2. Kinetic energy of unaffected part vs. surface energy
3. Drops cannot expel trapped gas film (bounce apart)
4. Drops form combined mass (coalesce)



$$We = \frac{\rho_L U^2 d_s}{\sigma}, \quad B = \frac{2b}{(d_s + d_l)}, \quad \Delta = \frac{d_s}{d_l}$$





Drop coalescence

Ashgriz & Poo, 1990

Grazing-coalescence boundary

Drops fly apart if rotational energy of colliding pair exceeds surface energy of combined pair

$$B = \sqrt{\frac{12}{5 We} \frac{(1 + \Delta^3)^{1/6}}{(1 + \Delta) \Delta^3} \left[1 + \Delta^2 - (1 + \Delta^3)^{2/3} \right]^{1/2}}$$

$0 < B < 1$
random number

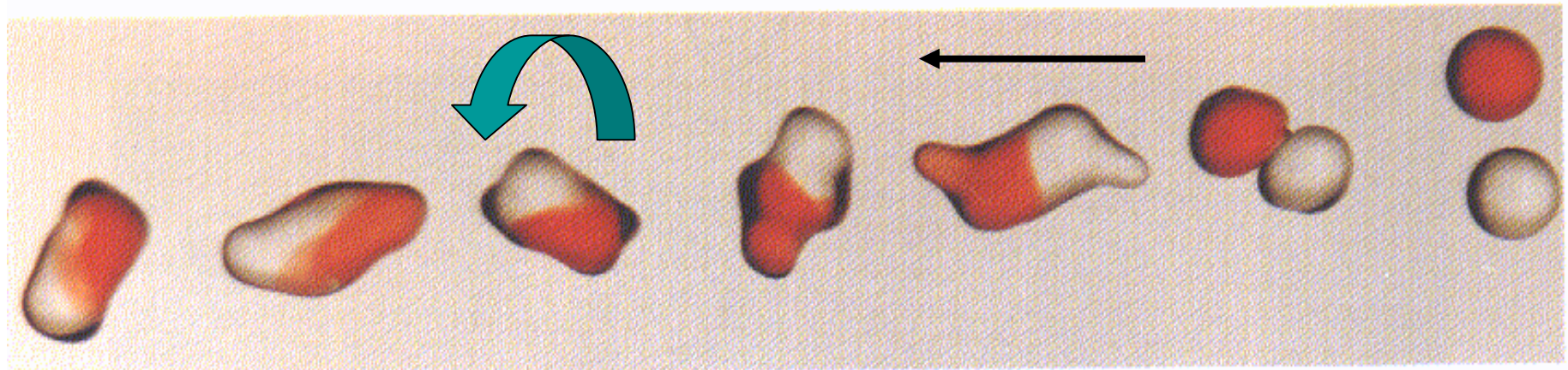


FIGURE 15. Coalescence collision at $\Delta = 1$, $We = 10$, and $B = 0.5$.





Grazing - Stretching Separation

Ashgriz & Poo, 1990

Energy and angular momentum conservation:

Grazing – drops move in same direction but at reduced velocity

Coalescence – mass average properties of colliding drops

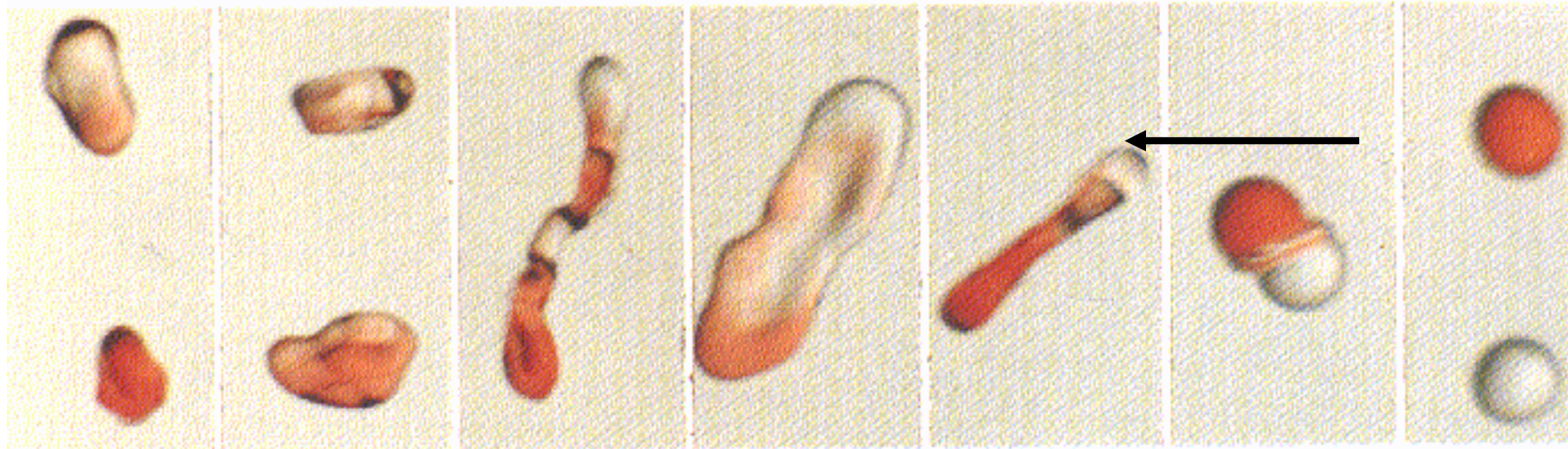


FIGURE 12. Stretching separation at $\Delta = 1$, $We = 53$, and $B = 0.38$.



Reflexive Separation

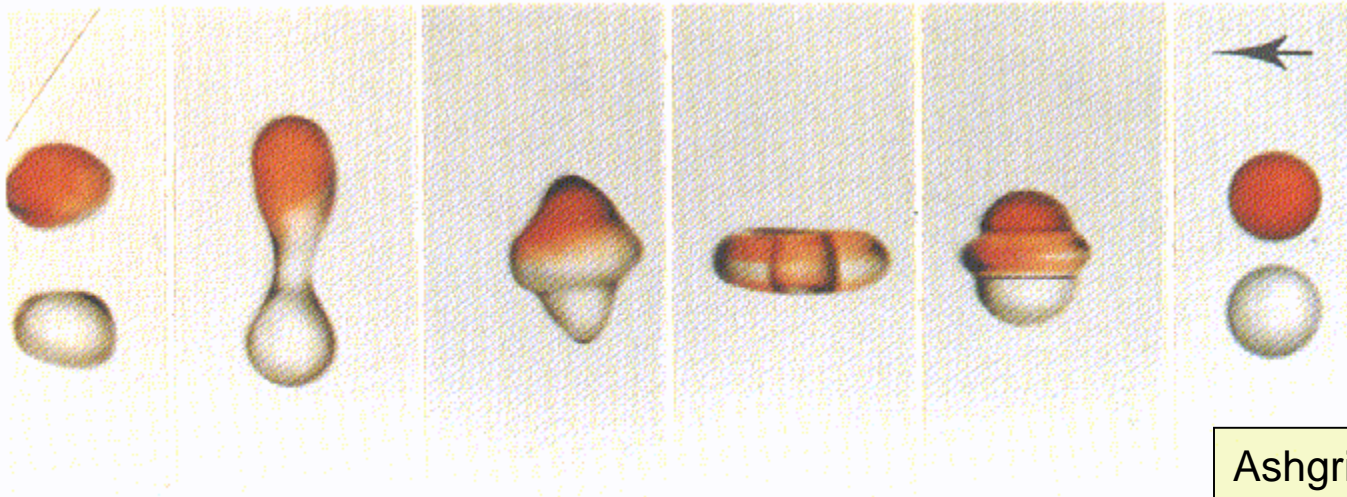
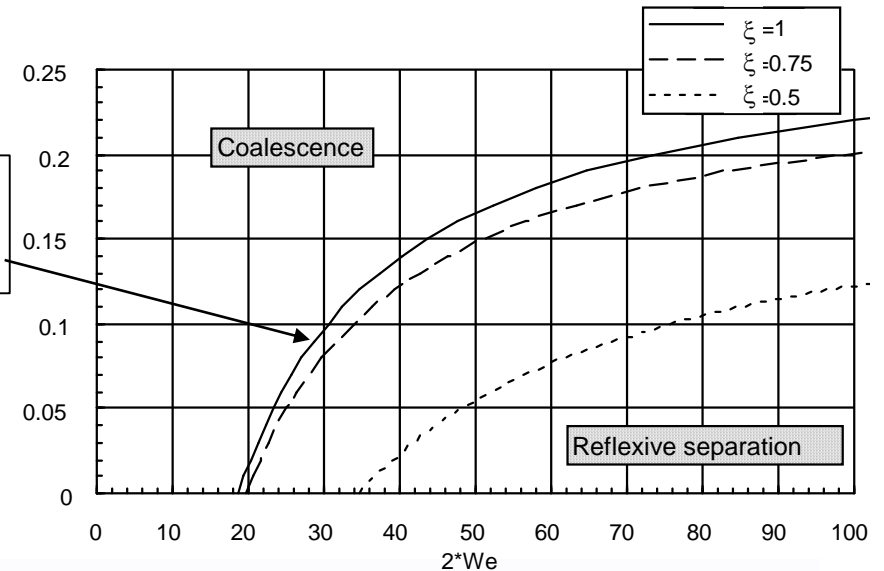
Tennison, SAE 980810

$$\frac{2 We}{\Delta (1 + \Delta^3)^2} (\Delta^6 \eta_1 + \eta_2) + 3 \left[4(1 + \Delta^2) - 7(1 + \Delta^3)^{2/3} \right] \geq 0$$

$$\eta_1 = 2(1 - \xi)^2 (1 - \xi^2)^{1/2} - 1$$

$$\eta_2 = 2(\Delta - \xi)^2 (\Delta^2 - \xi^2)^{1/2} - \Delta^3$$

with $\xi = B(1 + \Delta) / 2$



Ashgriz & Poo, 1990

FIGURE 5. Reflexive separation with no satellite for $\Delta = 1$, $We = 23$, and $B = 0.05$.





Summary

The Lagrangian Drop/Eulerian Fluid (LDEF) Discrete Drop model is the work-horse approach in commercial codes for simulating 2-phase flows.

Detailed models are available for use in engine CFD models to describe the effects of injector nozzle flow, and liquid and gas properties on spray formation and drop breakup physics.

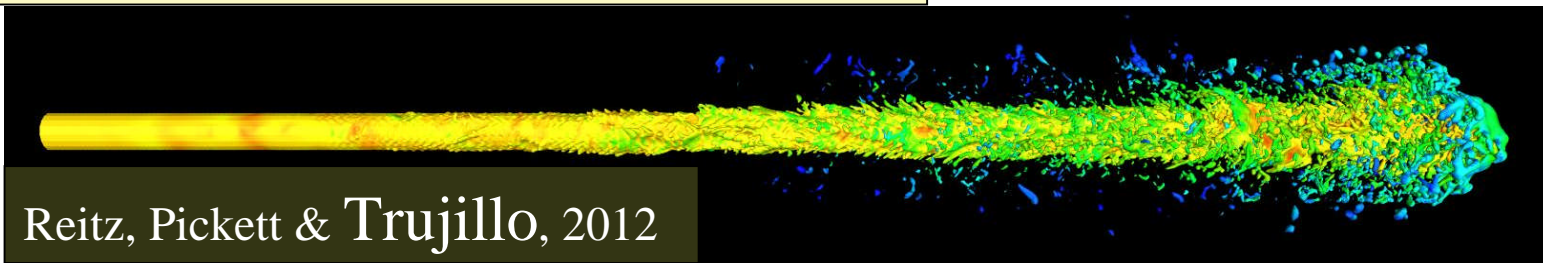
Due to the importance of sprays in applications, research is still needed. Recent experimental and modeling work can be accessed through ILASS and ICLASS conference papers and the Atomization and Sprays journal.

Significant progress is being made using LES/DNS spray modeling with high resolution experimental diagnostics to validate engine CFD spray models.

Ballistic imaging: Linne, 2009; **X-Ray imaging:** Liu SAE paper 2010-01-0877

LES: Villiers & Gosman, LES Primary Diesel Spray Atomization, SAE 2004-01-0100

DNS: Near field spray modeling (Trujillo - ERC)



Reitz, Pickett & Trujillo, 2012

